# Superadditivity of the convex closure of the output entropy of a quantum channel 

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One of the main recent achievements of quantum information theory [1] is the proof of the equivalence of several (sub-, super-) additivity conjectures for finite-dimensional quantum channels and systems [2] (see the survey in [3]). The main purpose of this paper is to give a generalization of this result to the infinite-dimensional case.

Let $\mathscr{H}$ and $\mathscr{H}^{\prime}$ be separable Hilbert spaces. A quantum channel is a linear completely positive trace-preserving map $\Phi: \mathfrak{T}(\mathscr{H}) \mapsto \mathfrak{T}\left(\mathscr{H}^{\prime}\right)$, where $\mathfrak{T}(\mathscr{H})$ is the ideal of all trace class operators on $\mathscr{H}$. In particular, $\Phi$ generates an affine map of the convex set $\mathfrak{S}(\mathscr{H})$ of states (that is, density operators) on the space $\mathscr{H}$ to the set $\mathfrak{S}\left(\mathscr{H}^{\prime}\right)$ of states on the space $\mathscr{H}^{\prime}[1]$.

Important characteristics of a quantum channel are the output entropy $H_{\Phi}(\rho)=$ $H(\Phi(\rho))$, a lower-semicontinuous concave function on the input state space with values in $[0,+\infty]$, and its convex closure (see [4]), denoted by $\widehat{H}_{\Phi}(\rho)$ and called the $\widehat{H}$-function of the channel $\Phi$. It was shown in [5] that the convex closure of the output entropy of an arbitrary quantum channel $\Phi$ is given by the expression

$$
\begin{equation*}
\widehat{H}_{\Phi}(\rho)=\inf _{\mu} \int_{\mathfrak{S}(\mathscr{H})} H_{\Phi}(\sigma) \mu(d \sigma), \tag{1}
\end{equation*}
$$

where the infimum is taken over all probability measures on $\mathfrak{S}(\mathscr{H})$ with barycentre $\rho$, and that this infimum is always achieved at some measure supported on the set of pure states.

A continuity condition for the $\widehat{H}$-function of an infinite-dimensional channel $\Phi$ was obtained in [5] (Proposition 7). This condition is equivalent to the condition

$$
\begin{equation*}
\left\{\lim _{n \rightarrow+\infty} H_{\Phi}\left(\rho_{n}\right)=H_{\Phi}\left(\rho_{0}\right)<+\infty\right\} \Longrightarrow\left\{\lim _{n \rightarrow+\infty} \widehat{H}_{\Phi}\left(\rho_{n}\right)=\widehat{H}_{\Phi}\left(\rho_{0}\right)\right\} . \tag{2}
\end{equation*}
$$

The superadditivity property of the $\widehat{H}$-function for channels $\Phi$ and $\Psi$ is the validity of the inequality

$$
\begin{equation*}
\widehat{H}_{\Phi \otimes \Psi}(\omega) \geqslant \widehat{H}_{\Phi}\left(\omega^{\mathscr{H}}\right)+\widehat{H}_{\Psi}\left(\omega^{\mathscr{K}}\right) \tag{3}
\end{equation*}
$$

for all $\omega \in \mathfrak{S}(\mathscr{H} \otimes \mathscr{K})$, where $\omega^{\mathscr{H}}=\operatorname{Tr}_{\mathscr{K}} \omega$ and $\omega^{\mathscr{K}}=\operatorname{Tr}_{\mathscr{H}} \omega$. This property implies the additivity of the minimal output entropy for the channels $\Phi$ and $\Psi$ (see [1], [3]). If $\Phi$ and $\Psi$ are partial traces, then the superadditivity of the $\widehat{H}$-function means the superadditivity of the Entanglement of Formation (EoF), which is an important characteristic of the state of a bipartite quantum system indicating the degree of entanglement of this state (see [1], [5]).

For finite-dimensional channels $\Phi$ and $\Psi$ the superadditivity of the $\widehat{H}$-function is equivalent to the additivity of the Holevo capacity of these channels with arbitrary constraints [6]. One of the obstacles to proving an analogous assertion for infinite-dimensional channels is the existence of superentangled states - pure states of a bipartite system having partial

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traces with infinite entropy (see Remark 4 in [7]). Namely, the existence of such states prevented, until now, proving the superadditivity of the $\widehat{H}$-function (and even the additivity of the minimal output entropy) for several classes of infinite-dimensional channels for which the additivity of the Holevo capacity with arbitrary constraints was derived in [7] from the corresponding finite-dimensional results [6], [8].

The problem of superentangled states can be partially solved by the following assertion based on the continuity condition (2) and on some other results in [5].

Lemma. Let $\Phi: \mathfrak{S}(\mathscr{H}) \mapsto \mathfrak{S}\left(\mathscr{H}^{\prime}\right)$ and $\Psi: \mathfrak{S}(\mathscr{K}) \mapsto \mathfrak{S}\left(\mathscr{K}^{\prime}\right)$ be arbitrary quantum channels. The inequality (3) holds for all $\omega$ in $\mathfrak{S}(\mathscr{H} \otimes \mathscr{K})$ if for an arbitrary finite-rank state $\omega_{0}$ in $\mathfrak{S}(\mathscr{H} \otimes \mathscr{K})$ with finite output entropy $H_{\Phi \otimes \Psi}\left(\omega_{0}\right)$ there exists a sequence $\left\{\omega_{n}\right\}$ of states in $\mathfrak{S}(\mathscr{H} \otimes \mathscr{K})$ such that the following properties are satisfied:
(i) $\lim _{n \rightarrow+\infty} \omega_{n}=\omega_{0}$ and $\lim _{n \rightarrow+\infty} H_{\Phi \otimes \Psi}\left(\omega_{n}\right)=H_{\Phi \otimes \Psi}\left(\omega_{0}\right)$;
(ii) the inequality (3) holds with $\omega=\omega_{n}$ for all $n \in \mathbb{N}$.

Using this lemma, Proposition 7 and Theorem 2 in [7], as well as Theorem 1 in [9], we obtain the following infinite-dimensional version of the results in [9], [6], [8].

Proposition. Let $\Psi$ be an arbitrary channel. Superadditivity of the $\widehat{H}$-function holds in the following cases: (i) $\Phi$ is a noiseless channel; (ii) $\Phi$ is an entanglement-breaking channel (see [3]); (iii) $\Phi$ is a channel complementary (see [9]) to an entanglement-breaking channel; (iv) $\Phi$ is a direct sum mixture (see [6]) of a noiseless channel and a channel $\Phi_{0}$ such that superadditivity of the $\widehat{H}$-function holds for the channels $\Phi_{0}$ and $\Psi$.

The above lemma makes it possible to prove the following assertion.
Theorem 1. If superadditivity of the $\widehat{H}$-function holds for all finite-dimensional quantum channels, then this property holds also for all infinite-dimensional quantum channels.

Corollary 1. If superadditivity of EoF holds for all states of a finite-dimensional bipartite quantum system, then this property holds also for all states of an infinite-dimensional bipartite quantum system.

Corollary 2. If additivity of the minimal output entropy holds for all finite-dimensional quantum channels, then this property holds also for all infinite-dimensional quantum channels.

Theorem 1 and Theorem 3 in [7] provide the following infinite-dimensional generalization of Shor's theorem [2].

Theorem 2. The following properties are equivalent: (i) additivity of the Holevo capacity holds for all infinite-dimensional quantum channels with arbitrary constraints; (ii) additivity of the minimal output entropy holds for all infinite-dimensional quantum channels; (iii) superadditivity of EoF holds for all states of an arbitrary infinite-dimensional bipartite quantum system.

The last property in Theorem 2 is equivalent to the superadditivity of the $\widehat{H}$-function for all infinite-dimensional quantum channels.

All the above assertions are proved in [10].

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